

Nonlinear free vibration analysis of levy plates using weak-form variational principle in polynomial displacement functions

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Abstract. This paper evaluates nonlinear free vibrations of Levy plates using Weak-Form variational principle in algebraic polynomial displacement functions. The energy functional of the plate problem was formulated using Weak-Form variational technique on the integral function of the Von Karman thin plate differential equations. The displacement functions were developed based on static deflection configurations of orthogonal beam network. The process of repeated direct integration on compatibility equation was used to determine the algebraic expressions for stress function. The amplitude of deflection which directly influences the geometric nonlinearity of the plate was determined using integration process on energy functional based on static equilibrium equations. The modal combination method was used to develop the stiffness and mass matrices respectively from the expressions of energy functional based on dynamic equilibrium equations. The numerical values of amplitudes of deflection at various aspect ratios were computed. Also, the first four nonlinear natural frequencies at various aspect ratios were numerically computed. The validation of the present study's results using the results from previous work found in literature shows satisfactory convergence, with an absolute mean error of 0.186 %. Conclusively, the application of Weak-Form variational principle in polynomial displacement functions provides satisfactory approximation to nonlinear dynamic analysis of Levy plates.

Keywords: algebraic polynomial, amplitude of deflection, energy functional, levy plate, nonlinear natural frequency, weak-form, variational principle.

1. Introduction

A levy plate in the context of this paper refers to any rectangular plate in which at least two opposite edge conditions are simply supported; and any of the two remaining edge conditions can either be clamped, simply supported or free. Engineering applications of levy plates are widely found in bridges, houses, marine structures, vehicular structures and airspace structures. In practice, almost all the rectangular plates engaged in the construction of these structures are classified as thin-walled structural components. Sequel to their thinness, the plates are often prone to large deformation in their engineering applications due to the influence of external disturbances. Consequently, the appropriate engineering analysis that is capable of simulating the actual response of thin rectangular plates to external disturbances must be the one formulated based on nonlinear mathematical model. However, the implication of using nonlinear mathematical model to determine the plate response is that the resulting nonlinear differential equations of equilibrium are likely going to impose complex mathematical challenges, and sometimes the closed-form solutions are absent. This problem is readily observed in the case of plates with mixed boundary conditions (majority of levy plates), which exhibit discontinuities in the edge supports [1].

For the fact that the resulting nonlinear differential equations from nonlinear method of analysis are tedious to tackle, approximation mathematical models have been developed to ameliorate the problem. Consequently, many of the investigators on large amplitude of vibrations of rectangular thin plates apply different approximation techniques to tackle the problem of

nonlinear vibrations of thin rectangular plates.

1.1. Related works

Engineering analysis of thin rectangular plates based on large amplitude of vibrations has for decade now posed interesting and challenging research fields. A number of approximation techniques have been developed for the purposes of investigating into large amplitude of vibrations of thin rectangular plates. [2] used Herrmann equations of motion in Galerkin technique to study the effect of initially applied stress on the nonlinear vibrations of flat plate. Finite element methods were used by [3-5] to investigate into large amplitude of vibrations of rectangular plates. Analytical technique on dynamic analogue of Von Karman differential equations was employed by [6] to investigate the effect of initial geometric imperfections and in-plane boundary conditions on the large amplitude of vibrations of angle- and cross-ply rectangular thin plates. [7] used levy solutions to study the vibration behaviour of multi-span rectangular plates. A combination of static and dynamic method-problem solving was used by [8] to investigate geometrically nonlinear free vibration of square plates with various boundary conditions. Second order asymptotic method was employed by [9] to investigate the effect of geometric imperfections on the amplitude of nonlinear vibrations of thin rectangular plates that were parametrically excited. A multiple scale method was used by [10] to investigate the nonlinear vibration of rectangular plates. Mesh-free least squares-based finite difference method was used by [11] to investigate large amplitude of vibrations of arbitrarily shaped thin plates. [12] used first order finite element theory to study free vibration of composite plates. A combination of Hamilton principle and Galerkin method was used by [13] to investigate nonlinear dynamic response of a simply supported functionally graded rectangular plate under the time-dependent thermalmechanical loads. Also, [14] used finite strip formulation to study nonlinear free vibration of plates. Furthermore, [15] used analytical technique, numerical and experimental analysis to study vibration behaviour of laminated composite plates under thermo-mechanical loading. The Galerkin-Vlasov method was used by [16] to obtain solution of free harmonic vibration of simply supported Kirchhoff plate.

1.2. Statement of problem and objectives

The energy methods have been proven to be very powerful for solving structural problems, especially those ones subjected to large deformations and large amplitude of vibrations. In analytical energy approach, the commonest approximation techniques appearing in most technical literature are the Hamilton's energy principle, the classical direct variational principle in Ritz method, and the residual energy method in Galerkin approach. Also, the most commonly used displacement functions are the trigonometric functions, either in the form of single or double Fourier series. There is need therefore, to explore other simple but accurate approximation techniques to be used in formulating energy functional for structural problems.

The primary objective of this study is to apply Weak-Form variational principle in formulating energy functional for nonlinear rectangular plate problem. The secondary objective is to develop and apply algebraic polynomial shape functions admissible to various levy plate boundary conditions.

2. Theoretical formulations

Fig. 1 shows a rectangular thin isotropic plate of levy type in which the edges $x = 0$ and $x = a$ are parallel and simply supported.

The boundary conditions of the edges $y = 0$ and $y = b$ are arbitrary; and each of them can either be clamped, simply supported or free. The plate is invariant in engineering properties of Young modulus E and Poisson's ratio μ , and it has constant mass density ρ . The plate lateral dimensions are defined by a and b as the length and width respectively.

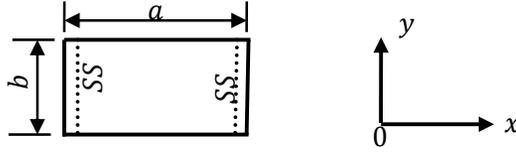


Fig. 1. A rectangular plate showing at least two edges parallel and simply supported

2.1. Formulation of energy functional

The Weak-form variational principle is used to formulate the energy functional of the plate problems. The plate is assumed to be thin; and it satisfies the integral function of Von Karman’s thin plate differential equations. The analogous Von Karman’s thin plate differential equations of motion are as expressed in Eqs. (1-2):

$$\nabla^4 w = \frac{1}{D} \left[q + \frac{\partial^2 F}{\partial y^2} \frac{\partial^2 w}{\partial x^2} - 2 \frac{\partial^2 F}{\partial x \partial y} \frac{\partial^2 w}{\partial x \partial y} + \frac{\partial^2 F}{\partial x^2} \frac{\partial^2 w}{\partial y^2} - \rho t \ddot{w} \right], \tag{1}$$

$$\nabla^4 F = Et \left[\left(\frac{\partial^2 w}{\partial x \partial y} \right)^2 - \frac{\partial^2 w}{\partial x^2} \frac{\partial^2 w}{\partial y^2} \right], \tag{2}$$

where E is Young modulus of plate materials; w is displacement functions; F is Airy’s stress functions; t is plate thickness; ρ is mass density of plate materials; q is uniformly applied lateral load; \ddot{w} is second derivative of w with respect to time T ; D is flexural rigidity of the plate

Let $V(x, y)$ be defined as weighting function which is admissible to boundary conditions of the levy plates. Then the variational statement is applied on the integral functions of Eq. (1) as expressed in Eq. (3):

$$0 = \int_0^a \int_0^b V(x, y) \left[\frac{\partial^4 w}{\partial x^4} + 2 \frac{\partial^4 w}{\partial x^2 \partial y^2} + \frac{\partial^4 w}{\partial y^4} - \frac{1}{D} \left(q + \frac{\partial^2 F}{\partial y^2} \frac{\partial^2 w}{\partial x^2} - 2 \frac{\partial^2 F}{\partial x \partial y} \frac{\partial^2 w}{\partial x \partial y} + \frac{\partial^2 F}{\partial x^2} \frac{\partial^2 w}{\partial y^2} - \rho t \ddot{w} \right) \right] dx dy, \tag{3}$$

where D is defined as expressed in Eq. (4):

$$D = \frac{Et^3}{12(1 - \mu^2)}, \tag{4}$$

where μ is Poisson’s ratio.

Let K be the reciprocal of D as given by Eq. (5):

$$K = \frac{1}{D}. \tag{5}$$

Integration by parts is exercised on Eq. (3) to trade differentiation from w and F to $V(x, y)$, and after simplification yields the expressions as given in Eq. (6):

$$\begin{aligned}
 0 = & \int_0^a \int_0^b \left[\frac{\partial^2 V}{\partial x^2} \frac{\partial^2 w}{\partial x^2} + 2 \frac{\partial^2 V}{\partial x \partial y} \frac{\partial^2 w}{\partial x \partial y} + \frac{\partial^2 V}{\partial y^2} \frac{\partial^2 w}{\partial y^2} \pm 2KF \frac{\partial^2 V}{\partial x \partial y} \frac{\partial^2 w}{\partial x \partial y} \pm KF \frac{\partial^2 V}{\partial y^2} \frac{\partial^2 w}{\partial x^2} \right. \\
 & \left. \pm KF \frac{\partial^2 V}{\partial x^2} \frac{\partial^2 w}{\partial y^2} - KVq + KVpt\dot{w} \right] dx dy \\
 & + \int_0^a \left\{ \left[2 \frac{\partial V}{\partial x} (M_{xy}) - V \left(\frac{\partial M_y}{\partial y} \right) + \frac{\partial V}{\partial y} (M_y) + 2KF \frac{\partial V}{\partial x} (M_{xy}) + KV \frac{\partial F}{\partial y} (M_x) \right. \right. \\
 & \left. \left. + 2KVF \left(\frac{\partial M_{xy}}{\partial x} \right) - KF \frac{\partial V}{\partial y} (M_x) \right]_0^b \right\} dx \\
 & + \int_0^b \left\{ \left[-V \left(\frac{\partial M_x}{\partial x} \right) + \frac{\partial V}{\partial x} (M_x) - 2V \left(\frac{\partial M_{xy}}{\partial y} \right) - KF \frac{\partial V}{\partial y} (M_{xy}) - KVF \left(\frac{\partial M_{xy}}{\partial y} \right) \right. \right. \\
 & \left. \left. - 2KV \frac{\partial F}{\partial y} (M_{xy}) - KF \frac{\partial V}{\partial x} (M_y) + KV \frac{\partial F}{\partial x} (M_y) \right]_0^a \right\} dy,
 \end{aligned} \tag{6}$$

where:

$$V = V(x, y). \tag{7}$$

The expressions obtained in Eq. (6) contain three identically satisfied expressions: the functional of the plate problem, which is the integrand as contained in the squared bracket; and two natural boundary conditions, which are the integrands as contained in the curl brackets. These natural boundary conditions are identically satisfied at the boundary of the plate, and as such they are dropped in the further analysis in this work.

By assuming that the plate vibrates harmonically and performs sinusoidal time response, the inertia term is simplified, and the expression for the functional is as expressed in Eq. (8):

$$\begin{aligned}
 0 = & \int_0^a \int_0^b \left[\frac{\partial^2 V}{\partial x^2} \frac{\partial^2 w}{\partial x^2} + 2 \frac{\partial^2 V}{\partial x \partial y} \frac{\partial^2 w}{\partial x \partial y} + \frac{\partial^2 V}{\partial y^2} \frac{\partial^2 w}{\partial y^2} \pm 2KF \frac{\partial^2 V}{\partial x \partial y} \frac{\partial^2 w}{\partial x \partial y} \pm KF \frac{\partial^2 V}{\partial y^2} \frac{\partial^2 w}{\partial x^2} \right. \\
 & \left. \pm KF \frac{\partial^2 V}{\partial x^2} \frac{\partial^2 w}{\partial y^2} - KVq - KVpt\omega^2 w \right] dx dy.
 \end{aligned} \tag{8}$$

Furthermore, for a free vibration analysis, the lateral load q is disregarded so that Eq. (8) is simplified as expressed in Eq. (9):

$$\begin{aligned}
 0 = & \int_0^a \int_0^b \left[\frac{\partial^2 V}{\partial x^2} \frac{\partial^2 w}{\partial x^2} + 2 \frac{\partial^2 V}{\partial x \partial y} \frac{\partial^2 w}{\partial x \partial y} + \frac{\partial^2 V}{\partial y^2} \frac{\partial^2 w}{\partial y^2} \pm 2KF \frac{\partial^2 V}{\partial x \partial y} \frac{\partial^2 w}{\partial x \partial y} \pm KF \frac{\partial^2 V}{\partial y^2} \frac{\partial^2 w}{\partial x^2} \right. \\
 & \left. \pm KF \frac{\partial^2 V}{\partial x^2} \frac{\partial^2 w}{\partial y^2} - \omega^2 MKVw \right] dx dy,
 \end{aligned} \tag{9}$$

where:

$$M = \rho * t. \tag{10}$$

The expression of Eq. (9) is a generic functional for dynamics of rectangular thin plates.

2.2. The boundary conditions

The analysis of geometrically nonlinear thin rectangular plates requires the consideration of both the out-of-plane and in-plane boundary conditions. In this study, the levy plates of simply supported-clamped-simply supported-clamped (SCSC), simply supported-clamped-simply supported-simply supported (SCSS), simply supported-simply supported-simply supported-simply supported (SSSS), simply supported-clamped-simply supported-free (SCSF), and simply supported-simply supported-simply supported-free (SSSF) rectangular plates respectively are considered. The general boundary conditions with respect to these rectangular plates are stated as given in Eq. (11-23) [17-19]:

The SCSC Rectangular plate:

$$u = v = w = w_0 = N_{xy} = M_x = M_{xy} = \frac{\partial^2 w_0}{\partial x^2} = 0, \quad x = 0, a, \tag{11}$$

$$u = v = w = w_0 = \frac{\partial w}{\partial y} = \frac{\partial w_0}{\partial y} = 0, \quad y = 0, b. \tag{12}$$

The SCSS rectangular plate:

$$u = v = w = w_0 = N_{xy} = M_x = M_{xy} = \frac{\partial^2 w_0}{\partial x^2} = 0, \quad x = 0, a, \tag{13}$$

$$u = v = w = w_0 = \frac{\partial w}{\partial y} = \frac{\partial w_0}{\partial y} = 0, \quad y = 0, \tag{14}$$

$$u = v = w = w_0 = N_{yx} = M_y = M_{yx} = \frac{\partial^2 w_0}{\partial y^2} = 0, \quad y = b. \tag{15}$$

The SSSS rectangular plate:

$$u = v = w = w_0 = N_{xy} = M_x = M_{xy} = \frac{\partial^2 w_0}{\partial x^2} = 0, \quad x = 0, a, \tag{16}$$

$$u = v = w = w_0 = N_{yx} = M_y = M_{yx} = \frac{\partial^2 w_0}{\partial y^2} = 0, \quad y = 0, b. \tag{17}$$

The SCSF rectangular plate:

$$u = v = w = w_0 = N_{xy} = M_x = M_{xy} = \frac{\partial^2 w_0}{\partial x^2} = 0, \quad x = 0, a, \tag{18}$$

$$u = v = w = w_0 = \frac{\partial w}{\partial y} = \frac{\partial w_0}{\partial y} = 0, \quad y = 0, \tag{19}$$

$$M_y = Q_y = N_y = N_{yx} = 0, \quad y = b. \tag{20}$$

The SSSF rectangular plate:

$$u = v = w = w_0 = N_{xy} = M_x = M_{xy} = \frac{\partial^2 w_0}{\partial x^2} = 0, \quad x = 0, a, \tag{21}$$

$$u = v = w = w_0 = \frac{\partial w}{\partial y} = \frac{\partial w_0}{\partial y} = 0, \quad y = 0, \tag{22}$$

$$M_y = Q_y = N_y = N_{yx} = 0, \quad y = b. \tag{23}$$

2.3. Determination of specially energy functional

From the generic energy functional of Eq. (9), the specially energy functional with respect to any set of rectangular boundary conditions is determined by applying the general boundary conditions of Eq. (11-23). Thus, the specially energy functional (SEF) for SCSC rectangular plate is given as expressed in Eq. (24):

$$0 = \int_0^a \int_0^b \left[\frac{\partial^2 V}{\partial x^2} \frac{\partial^2 w}{\partial x^2} + 2 \frac{\partial^2 V}{\partial x \partial y} \frac{\partial^2 w}{\partial x \partial y} + \frac{\partial^2 V}{\partial y^2} \frac{\partial^2 w}{\partial y^2} + KF \frac{\partial^2 V}{\partial x \partial y} \frac{\partial^2 w}{\partial x \partial y} + KF \frac{\partial^2 V}{\partial y^2} \frac{\partial^2 w}{\partial x^2} + KF \frac{\partial^2 V}{\partial x^2} \frac{\partial^2 w}{\partial y^2} - \omega^2 MKVw \right] dx dy. \quad (24)$$

The SEF for SCSS and SSSS rectangular plates is as given by Eq. (25):

$$0 = \int_0^a \int_0^b \left[\frac{\partial^2 V}{\partial x^2} \frac{\partial^2 w}{\partial x^2} + 2 \frac{\partial^2 V}{\partial x \partial y} \frac{\partial^2 w}{\partial x \partial y} + \frac{\partial^2 V}{\partial y^2} \frac{\partial^2 w}{\partial y^2} + KF \frac{\partial^2 V}{\partial y^2} \frac{\partial^2 w}{\partial x^2} + KF \frac{\partial^2 V}{\partial x^2} \frac{\partial^2 w}{\partial y^2} - \omega^2 MKVw \right] dx dy. \quad (25)$$

And the SEF for SCSF and SSSF rectangular plates is as given in Eq. (26):

$$0 = \int_0^a \int_0^b \left[\frac{\partial^2 V}{\partial x^2} \frac{\partial^2 w}{\partial x^2} + 2 \frac{\partial^2 V}{\partial x \partial y} \frac{\partial^2 w}{\partial x \partial y} + \frac{\partial^2 V}{\partial y^2} \frac{\partial^2 w}{\partial y^2} + KF \frac{\partial^2 V}{\partial y^2} \frac{\partial^2 w}{\partial x^2} - \omega^2 MKVw \right] dx dy. \quad (26)$$

2.4. Development of displacement functions

Fig. 2 shows a levy plate consisting of an idealized series of orthogonal beam network in x and y -directions. It is assumed that each beam element running along any coordinate direction is a good representative of other series of beam elements that run along the same direction.

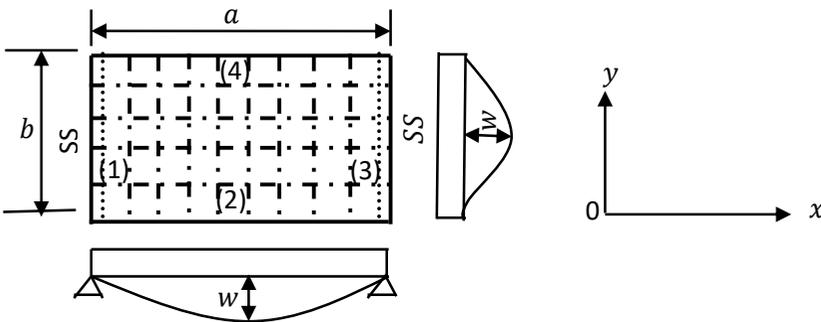


Fig. 2. A Levy plate showing a series of orthogonal beam network and deflection configurations

Let $w(x)$ and $w(y)$ denote the algebraic polynomial displacement functions in x and y -directions as expressed in Eq. (27) and Eq. (28) respectively:

$$w(x) = \sum_{i=0}^r c_i x^i, \quad (27)$$

$$w(y) = \sum_{j=0}^r d_j y^j, \tag{28}$$

where c_i and d_j are undetermined coefficients.

Numerically, it has been proven that a five-term polynomials provides necessary and sufficient satisfaction for completeness as shape functions for a beam element. Therefore, the value of r in the summation is four. The out-of-plane boundary conditions for clamped edge, simply supported edge and free edge are as expressed in Eq. (29), Eq. (30) and Eq. (31) respectively:

$$w(\cdot, \cdot) = w'(\cdot, \cdot) = 0, \tag{29}$$

$$w(\cdot, \cdot) = w''(\cdot, \cdot) = 0, \tag{30}$$

$$w'(\cdot, \cdot) = w''(\cdot, \cdot) = w'''(\cdot, \cdot) = 0, \tag{31}$$

where $w(\cdot, \cdot)$ represents deflection function, $w'(\cdot, \cdot)$, $w''(\cdot, \cdot)$ and $w'''(\cdot, \cdot)$ represent first, second and third derivatives with respect to the coordinates respectively.

Suppose Fig. 2 depicts a levy plate of SCSC rectangular plate, then the deflection along x -direction is expressed as given in Eq. (32):

$$w(x) = c_0 + c_1 x + c_2 x^2 + c_3 x^3 + c_4 x^4. \tag{32}$$

By applying the boundary conditions of Eq. (30) in Eq. (32), and subsequently carrying out simplification, gives Eq. (33):

$$w(x) = c_4 * a^4 \left[\left(\frac{x}{a}\right) - 2 \left(\frac{x}{a}\right)^3 + \left(\frac{x}{a}\right)^4 \right]. \tag{33}$$

However, to account for series of beam mode shapes running along y -direction, Eq. (33) is modified as shown in Eq. (34):

$$w(x) = c_4 * a^4 \left[\left(\frac{x}{a}\right)^m - 2 \left(\frac{x}{a}\right)^{m+2} + \left(\frac{x}{a}\right)^{m+3} \right]. \tag{34}$$

Similarly suppose the edges Eqs. (2) and (4) of Fig. 2 are clamped, then the polynomial expression is as expressed in Eq. (35):

$$w(y) = d_0 + d_1 y + d_2 y^2 + d_3 y^3 + d_4 y^4. \tag{35}$$

By applying the boundary conditions of Eq. (29) in Eq. (35) and after simplifying, yields Eq. (36):

$$w(y) = d_4 * b^4 \left[\left(\frac{y}{b}\right)^2 - 2 \left(\frac{y}{b}\right)^3 + \left(\frac{y}{b}\right)^4 \right]. \tag{36}$$

Also, to account for series of beam mode shapes running along x -direction, Eq. (36) is modified as shown in Eq. (37):

$$w(y) = d_4 * b^4 \left[\left(\frac{y}{b}\right)^{n+1} - 2 \left(\frac{y}{b}\right)^{n+2} + \left(\frac{y}{b}\right)^{n+3} \right], \tag{37}$$

where $m, n = 1, 2, 3, \dots$

The rectangular plate static mode shape is the product of Eq. (34) and Eq. (37) as expressed in Eq. (38):

$$w(x, y) = A_{mn} \left[\left(\frac{x}{a}\right)^m - 2 \left(\frac{x}{a}\right)^{m+2} + \left(\frac{x}{a}\right)^{m+3} \right] \left[\left(\frac{y}{b}\right)^{n+1} - 2 \left(\frac{y}{b}\right)^{n+2} + \left(\frac{y}{b}\right)^{n+3} \right], \quad (38)$$

where A_{mn} is as expressed in Eq. (39):

$$A_{mn} = (c_4 * a^4)(d_o * b^4). \quad (39)$$

The parameter A_{mn} is the amplitude of deflection function. Thus Eq. (38) is the static algebraic polynomial displacement function for SCSC rectangular plate. The same procedure is used to develop the static algebraic polynomial mode shapes for SCSS, SSSS, SCSF and SSSF rectangular plates; and they are as expressed in Eq. (40-43) respectively:

$$w(x, y) = A_{mn} \left[\left(\frac{x}{a}\right)^m - 2 \left(\frac{x}{a}\right)^{m+2} + \left(\frac{x}{a}\right)^{m+3} \right] \left[\frac{3}{2} \left(\frac{y}{b}\right)^{n+1} - \frac{5}{2} \left(\frac{y}{b}\right)^{n+2} + \left(\frac{y}{b}\right)^{n+3} \right], \quad (40)$$

$$w(x, y) = A_{mn} \left[\left(\frac{x}{a}\right)^m - 2 \left(\frac{x}{a}\right)^{m+2} + \left(\frac{x}{a}\right)^{m+3} \right] \left[\left(\frac{y}{b}\right)^n - 2 \left(\frac{y}{b}\right)^{n+2} + \left(\frac{y}{b}\right)^{n+3} \right], \quad (41)$$

$$w(x, y) = A_{mn} \left[\left(\frac{x}{a}\right)^m - 2 \left(\frac{x}{a}\right)^{m+2} + \left(\frac{x}{a}\right)^{m+3} \right] \left[6 \left(\frac{y}{b}\right)^{n+1} - 4 \left(\frac{y}{b}\right)^{n+2} + \left(\frac{y}{b}\right)^{n+3} \right], \quad (42)$$

$$w(x, y) = A_{mn} \left[\left(\frac{x}{a}\right)^m - 2 \left(\frac{x}{a}\right)^{m+2} + \left(\frac{x}{a}\right)^{m+3} \right] \left[2 \left(\frac{y}{b}\right)^n - 2 \left(\frac{y}{b}\right)^{n+2} + \left(\frac{y}{b}\right)^{n+3} \right]. \quad (43)$$

2.5. The stress function

One of the unknown parameters in the expression for the energy functional is the Airy's stress function F . In this work, direct integration process on compatibility Eq. (2) is used to determine the expression for stress function. Thus, a repeated integration process on the right hand side of Eq. (2) is performed four times with respect to x -coordinate; and then four times with respect to y -coordinate as expressed in Eq. (44):

$$F_{mn} = \int_x \int_x \int_x \int_x \int_y \int_y \int_y \int_y \left[\left(\frac{\partial^2 w(mn)}{\partial x \partial y} \right)^2 - \frac{\partial^2 w(mn)}{\partial x^2} \frac{\partial^2 w(mn)}{\partial y^2} \right], \quad (44)$$

where $\int_x = \int(\cdot)dx$ and $\int_y = \int(\cdot)dy$; and for $m = 1, 2$ and $n = 1, 2, 3$.

In this study, the following designations for describing mode shapes are adopted: for $m = n = 1$, then $w(1,1)$ denotes the first mode shape; for $m = 1$ and $n = 2$, then $w(1,2)$ denotes the second mode shape; for $m = 1$ and $n = 3$, then $w(1,3)$ denotes the third mode shape; and for $m = 2$ and $n = 2$, then $w(2,2)$ denotes the fourth mode shape. For instance, at mode 1, the mode shape with respect to SCSC rectangular plate is as given in Eq. (45):

$$w(1,1) = A_{11} \left[\left(\frac{x}{a}\right)^1 - 2 \left(\frac{x}{a}\right)^3 + \left(\frac{x}{a}\right)^4 \right] \left[\left(\frac{y}{b}\right)^2 - 2 \left(\frac{y}{b}\right)^3 + \left(\frac{y}{b}\right)^4 \right]. \quad (45)$$

Therefore, the corresponding algebraic stress function F_{11} with respect to SCSC rectangular plate evaluated using Eq. (44) is as expressed in Eq. (46):

$$A_{11} = \frac{-1.25992\xi_1}{\left[Z_1 + \sqrt{Z_2 + Z_3\xi_1^3}\right]^{\frac{1}{3}}} + \zeta_1 \left\{ 1.99059 * 10^7 \left[Z_1 + \sqrt{Z_2 + Z_3\xi_1^3} \right]^{\frac{1}{3}} \right\}, \quad (50)$$

where:

$$\xi_1 = (0.0393651 + 0.0185034\beta^2 + 0.00761907\beta^4), \quad (51)$$

$$Z_1 = 3.17964 * 10^{-17} a^{12} K^3 q Q^2 \beta^{12}, \quad (52)$$

$$Z_2 = 1.01101 * 10^{-33} a^{24} K^6 Q^2 Q^4 \beta^{24}, \quad (53)$$

$$Z_3 = 2.53561 * 10^{-22} a^{12} K^3 Q^3 \beta^{12}, \quad (54)$$

$$\zeta_1 = \frac{1}{a^4 K Q \beta^4}. \quad (55)$$

Thus, the expressions for A_{12} , A_{13} , and A_{22} are determined using the same procedure.

2.7. Determination of stiffness and mass matrices

The vibration of each mode shape of a rectangular plate consists of a product of two perpendicular motions in x and y -coordinates respectively. During natural vibrations, there exist modal interactions among different modes. Consequently, in this work, the mode shape w and the weighting function V are interactively combined to effect these modal interactions. Thus, to determine the natural frequencies, the expressions representing the energy functional of Eq. (24-26) are invoked and modified in accordance with these modal interactions as expressed in Eq. (56-58) respectively:

$$\int_0^a \int_0^b \left[\frac{\partial^2 V}{\partial x^2} \frac{\partial^2 w}{\partial x^2} + 2 \frac{\partial^2 V}{\partial x \partial y} \frac{\partial^2 w}{\partial x \partial y} + \frac{\partial^2 V}{\partial y^2} \frac{\partial^2 w}{\partial y^2} + K F_V F_w \frac{\partial^2 V}{\partial x \partial y} \frac{\partial^2 w}{\partial x \partial y} + K F_V F_w \frac{\partial^2 V}{\partial y^2} \frac{\partial^2 w}{\partial x^2} + K F_V F_w \frac{\partial^2 V}{\partial x^2} \frac{\partial^2 w}{\partial y^2} \right] dx dy = \omega^2 M K \int_0^a \int_0^b V w dx dy, \quad (56)$$

$$\int_0^a \int_0^b \left[\frac{\partial^2 V}{\partial x^2} \frac{\partial^2 w}{\partial x^2} + 2 \frac{\partial^2 V}{\partial x \partial y} \frac{\partial^2 w}{\partial x \partial y} + \frac{\partial^2 V}{\partial y^2} \frac{\partial^2 w}{\partial y^2} + K F_V F_w \frac{\partial^2 V}{\partial y^2} \frac{\partial^2 w}{\partial x^2} + K F_V F_w \frac{\partial^2 V}{\partial x^2} \frac{\partial^2 w}{\partial y^2} \right] dx dy = \omega^2 M K \int_0^a \int_0^b V w dx dy, \quad (57)$$

$$\int_0^a \int_0^b \left[\frac{\partial^2 V}{\partial x^2} \frac{\partial^2 w}{\partial x^2} + 2 \frac{\partial^2 V}{\partial x \partial y} \frac{\partial^2 w}{\partial x \partial y} + \frac{\partial^2 V}{\partial y^2} \frac{\partial^2 w}{\partial y^2} + K F_V F_w \frac{\partial^2 V}{\partial y^2} \frac{\partial^2 w}{\partial x^2} \right] dx dy = \omega^2 M K \int_0^a \int_0^b V w dx dy, \quad (58)$$

where F_V and F_w are the stress functions associated with mode shapes corresponding to weighting function and displacement function respectively. For instance, to determine the stiffness component k_{11} at first mode, let $V = w = w(1,1)$; to determine the stiffness component $k_{21} = k_{12}$, let $V = w(1,2)$ and $w = w(1,1)$; to determine the stiffness component $k_{31} = k_{13}$, let $V = w(1,3)$ and $w = w(1,1)$; to determine the stiffness component $k_{41} = k_{14}$, let $V = w(2,2)$ and $w = w(1,1)$ and so on. For instance, the mode shapes for SCSC plate are as expressed in

Eqs. (59-62):

$$w(1,1) = A_{11} \left[\left(\frac{x}{a} \right) - 2 \left(\frac{x}{a} \right)^3 + \left(\frac{x}{a} \right)^4 \right] \left[\left(\frac{y}{b} \right)^2 - 2 \left(\frac{y}{b} \right)^3 + \left(\frac{y}{b} \right)^4 \right], \quad (59)$$

$$w(1,2) = A_{12} \left[\left(\frac{x}{a} \right) - 2 \left(\frac{x}{a} \right)^3 + \left(\frac{x}{a} \right)^4 \right] \left[\left(\frac{y}{b} \right)^3 - 2 \left(\frac{y}{b} \right)^4 + \left(\frac{y}{b} \right)^5 \right], \quad (60)$$

$$w(1,3) = A_{13} \left[\left(\frac{x}{a} \right) - 2 \left(\frac{x}{a} \right)^3 + \left(\frac{x}{a} \right)^4 \right] \left[\left(\frac{y}{b} \right)^4 - 2 \left(\frac{y}{b} \right)^5 + \left(\frac{y}{b} \right)^6 \right], \quad (61)$$

$$w(2,2) = A_{22} \left[\left(\frac{x}{a} \right)^2 - 2 \left(\frac{x}{a} \right)^4 + \left(\frac{x}{a} \right)^5 \right] \left[\left(\frac{y}{b} \right)^3 - 2 \left(\frac{y}{b} \right)^4 + \left(\frac{y}{b} \right)^5 \right]. \quad (62)$$

Then the stiffness matrix for SCSC rectangular plate is determined using the expressions given in Eqs. (63-67):

$$\int_0^a \int_0^b \left[\left(\frac{\partial^2 w(1,1)}{\partial x^2} \right)^2 + 2 \left(\frac{\partial^2 w(1,1)}{\partial x \partial y} \right)^2 + \left(\frac{\partial^2 w(1,1)}{\partial y^2} \right)^2 + K(F_{11})^2 \left(\frac{\partial^2 w(1,1)}{\partial x \partial y} \right)^2 + 2K(F_{11})^2 \left(\frac{\partial^2 w(1,1)}{\partial x^2} \right) \left(\frac{\partial^2 w(1,1)}{\partial y^2} \right) \right] dx dy = k_{11}, \quad (63)$$

$$\int_0^a \int_0^b \left[\frac{\partial^2 w(1,2)}{\partial x^2} \frac{\partial^2 w(1,1)}{\partial x^2} + 2 \frac{\partial^2 w(1,2)}{\partial x \partial y} \frac{\partial^2 w(1,1)}{\partial x \partial y} + \frac{\partial^2 w(1,2)}{\partial y^2} \frac{\partial^2 w(1,1)}{\partial y^2} + KF_{12}F_{11} \frac{\partial^2 w(1,2)}{\partial x \partial y} \frac{\partial^2 w(1,1)}{\partial x \partial y} + KF_{12}F_{11} \frac{\partial^2 w(1,2)}{\partial y^2} \frac{\partial^2 w(1,1)}{\partial x^2} + KF_{12}F_{11} \frac{\partial^2 w(1,2)}{\partial x^2} \frac{\partial^2 w(1,1)}{\partial y^2} \right] dx dy = k_{21} = k_{12}, \quad (64)$$

$$\int_0^a \int_0^b \left[\frac{\partial^2 w(1,3)}{\partial x^2} \frac{\partial^2 w(1,1)}{\partial x^2} + 2 \frac{\partial^2 w(1,3)}{\partial x \partial y} \frac{\partial^2 w(1,1)}{\partial x \partial y} + \frac{\partial^2 w(1,3)}{\partial y^2} \frac{\partial^2 w(1,1)}{\partial y^2} + KF_{13}F_{11} \frac{\partial^2 w(1,3)}{\partial x \partial y} \frac{\partial^2 w(1,1)}{\partial x \partial y} + KF_{13}F_{11} \frac{\partial^2 w(1,3)}{\partial y^2} \frac{\partial^2 w(1,1)}{\partial x^2} + KF_{13}F_{11} \frac{\partial^2 w(1,3)}{\partial x^2} \frac{\partial^2 w(1,1)}{\partial y^2} \right] dx dy = k_{31} = k_{13}, \quad (65)$$

$$\int_0^a \int_0^b \left[\frac{\partial^2 w(2,2)}{\partial x^2} \frac{\partial^2 w(1,1)}{\partial x^2} + 2 \frac{\partial^2 w(2,2)}{\partial x \partial y} \frac{\partial^2 w(1,1)}{\partial x \partial y} + \frac{\partial^2 w(2,2)}{\partial y^2} \frac{\partial^2 w(1,1)}{\partial y^2} + KF_{22}F_{11} \frac{\partial^2 w(2,2)}{\partial x \partial y} \frac{\partial^2 w(1,1)}{\partial x \partial y} + KF_{22}F_{11} \frac{\partial^2 w(2,2)}{\partial y^2} \frac{\partial^2 w(1,1)}{\partial x^2} + KF_{22}F_{11} \frac{\partial^2 w(2,2)}{\partial x^2} \frac{\partial^2 w(1,1)}{\partial y^2} \right] dx dy = k_{41} = k_{14}, \quad (66)$$

$$\int_0^a \int_0^b \left[\left(\frac{\partial^2 w(2,2)}{\partial x^2} \right)^2 + 2 \left(\frac{\partial^2 w(2,2)}{\partial x \partial y} \right)^2 + \left(\frac{\partial^2 w(2,2)}{\partial y^2} \right)^2 + K(F_{22})^2 \left(\frac{\partial^2 w(2,2)}{\partial x \partial y} \right)^2 + 2K(F_{22})^2 \left(\frac{\partial^2 w(2,2)}{\partial x^2} \right) \left(\frac{\partial^2 w(2,2)}{\partial y^2} \right) \right] dx dy = k_{44}. \quad (67)$$

The resulting stiffness matrix is given as expressed in Eq. (68):

$$[\mathbf{K}] = \begin{bmatrix} k_{11} & k_{12} & k_{13} & k_{14} \\ k_{21} & k_{22} & k_{23} & k_{24} \\ k_{31} & k_{32} & k_{33} & k_{34} \\ k_{41} & k_{42} & k_{43} & k_{44} \end{bmatrix}. \quad (68)$$

The same procedure is used to develop other required stiffness matrices for SCSS, SSSS, SCSF and SSSF respectively. In similar approach, the right hand side (RHS) of Eq. (56), Eq. (57) and Eq. (58) respectively will yield the mass matrices after due integration process. Thus, for SCSC rectangular plate, the mass matrix is determined by using expressions given in Eqs. (69-72):

$$\omega^2 MK \int_0^a \int_0^b [(w(1,1))^2] dx dy = b_{11}, \quad (69)$$

$$\omega^2 MK \int_0^a \int_0^b [w(1,2) * w(1,1)] dx dy = b_{21} = b_{12}, \quad (70)$$

$$\omega^2 MK \int_0^a \int_0^b [w(1,3) * w(1,1)] dx dy = b_{31} = b_{13}, \quad (71)$$

$$\omega^2 MK \int_0^a \int_0^b [(w(2,2))^2] dx dy = b_{44}. \quad (72)$$

The resulting mass matrix is as given in Eq. (73):

$$[\mathbf{B}] = \omega^2 \begin{bmatrix} b_{11} & b_{12} & b_{13} & b_{14} \\ b_{21} & b_{22} & b_{23} & b_{24} \\ b_{31} & b_{32} & b_{33} & b_{34} \\ b_{41} & b_{42} & b_{43} & b_{44} \end{bmatrix} MK. \quad (73)$$

The same approach is used to develop the mass matrices for SCSS, SSSS, SCSF and SSSF respectively.

3. Numerical evaluation and discussion of results

Natural vibrations of rectangular isotropic plates under geometrical nonlinearity are studied in this paper using weak-form variational principle. Numerical evaluations are required for adequate validation of the results of the present study with the results found in existing literature. The plate parameters used in this study are: $E = 10.92$ MPa, $\rho = 100$ kgm⁻³, $\mu = 0.3$, $a = 1.0$ m, $t = 0.01$ m, $b = \text{open}$, and $\beta = b/a$, which were previously used by [3, 11, 20, 21] in their separate studies using different approximation techniques. Thus, by using aspect ratio, $\beta = 1$, these parameters were appropriately substituted and the numerical values for the elements in the stiffness matrix of Eq. (68) are as given in Eq. (74):

$$[\mathbf{K}] = \begin{bmatrix} 0.0632071 & 0.0415130 & 0.0221544 & 0.0289220 \\ 0.0415130 & 0.0419203 & 0.0288422 & 0.0292578 \\ 0.0221544 & 0.0288422 & 0.0235024 & 0.0201416 \\ 0.0289220 & 0.0292578 & 0.0201416 & 0.0293874 \end{bmatrix}. \quad (74)$$

Similarly, by substituting these parameters appropriately, the numerical values for the elements in the mass matrix of Eq. (73) are as given in Eq. (75):

$$[B] = \omega^2 \begin{bmatrix} 7.46967 & 4.93073 & 2.62527 & 3.44405 \\ 4.93073 & 3.55067 & 2.02177 & 2.48009 \\ 2.62527 & 2.02177 & 1.21446 & 1.41218 \\ 3.44405 & 2.48009 & 1.41218 & 1.96091 \end{bmatrix} 10^{-5}. \tag{75}$$

In the matrix manipulation, suppose the difference between $[K]$ and $[B]$ is $[A]$, then:

$$[A] = [K] - [B]. \tag{76}$$

The determinant of $[A]$ is thus:

$$\det([A]) = \det([K] - [B]). \tag{77}$$

For nontrivial solution of the eigen-function, Eq. (77) is set equal to zero, as expressed in Eq. (78):

$$\det([K] - [B]) = 0. \tag{78}$$

The expression of Eq. (78) can be simplified numerically by pre-multiplying Eq. (78) by the inverse of Eq. (75) as expressed in Eq. (79):

$$\det([B]^{-1}[K] - [B]^{-1}[B]) = 0. \tag{79}$$

The result of the manipulation of Eq. (79) yields an expression as given in Eq. (80):

$$\det \begin{bmatrix} (1323.27 - R) & -2441.51 & 2584.17 & -1723.73 \\ -1482.14 & (4108.7 - R) & -13388.7 & 177.234 \\ 1468.96 & 824.251 & (18640.5 - R) & 536.425 \\ -32.5375 & -9.94004 & -2.26802 & (3915.67 - R) \end{bmatrix} = 0, \tag{80}$$

where R is $R = \omega^2$.

The solution of the resulting quartic algebraic equation in terms of R and subsequently in terms of ω gives the nonlinear natural frequencies of vibration of the rectangular SCSC plate at aspect ratio equal to one. Therefore, by carrying out the same procedure with respect to each of the SCSS, SSSS, SCSF and SSSF rectangular plates respectively yields the desired natural frequencies for the respective rectangular plate up to first four lowest natural frequencies.

However, Table 1 presents the numerical values of the variation of the amplitudes of plate deflection with aspect ratios for the SCSC, SCSS, SSSS, SCSF and SSSF rectangular plates. The amplitudes of deflection are the measures of the rigidity characteristics of rectangular plates; and under geometrical nonlinearity of rectangular plates, these geometric characteristics shift from rigid to flexible plate conditions [21]. The amplitudes of deflection are strong parameters that affect the nonlinear flexural frequencies of rectangular plates based on large deflection.

From Table 1, it was observed that the amplitudes of deflection for the SCSC, SCSS, SSSS and SSSF rectangular plates exhibit hard-spring type with the aspect ratios; but with the exception of the amplitudes of deflection of the third mode of the SSSF rectangular plate boundary conditions, which exhibit soft-spring type. On the other hand, it was observed that the amplitudes of deflection for the SCSF rectangular plate boundary conditions exhibit entirely soft-spring type with the aspect ratios.

Also, Table 2 presents the numerical values of nonlinear natural frequencies for the SCSC, SCSS, SSSS, SCSF and SSSF rectangular plates based on the above given plate parameters at various aspect ratios.

Table 1. Variation of amplitudes of deflection with aspect ratios

Boundary condition	Mode	Amplitude of deflection, A_{mn}					
		Aspect ratio, β					
		$\beta = 0.25$	$\beta = 0.50$	$\beta = 0.75$	$\beta = 1.0$	$\beta = 1.25$	$\beta = 1.50$
SCSC	A_{11}	0.0063	0.0919	0.3960	0.9779	1.6535	2.1545
	A_{12}	0.0074	0.1102	0.4925	1.2911	2.3935	3.4001
	A_{13}	0.0064	0.0974	0.4527	1.2603	2.5691	4.1174
	A_{22}	0.0130	0.1863	0.7711	1.8036	3.0012	4.0434
SCSS	A_{11}	0.0061	0.0824	0.3160	0.6845	1.0382	1.2800
	A_{12}	0.0054	0.0760	0.3106	0.7364	1.2454	1.6731
	A_{13}	0.0032	0.0477	0.2116	0.5574	1.0770	1.6642
	A_{22}	0.0094	0.1242	0.4556	0.9455	1.4380	1.8363
SSSS	A_{11}	0.0057	0.0664	0.2121	0.3788	0.4955	0.5639
	A_{12}	0.0034	0.0466	0.1837	0.4140	0.6552	0.8314
	A_{13}	0.0020	0.0287	0.1253	0.3220	0.6009	0.8879
	A_{22}	0.0059	0.0750	0.2604	0.5091	0.7352	0.9036
SCSF	A_{11}	11.4724	4.2488	2.8316	2.3019	2.0428	1.8960
	A_{12}	4.8839	1.6481	1.0031	0.7499	0.6163	0.5327
	A_{13}	7.1802	2.1422	1.1845	0.8296	0.6509	0.5432
	A_{22}	9.8596	3.6792	2.4892	2.0541	1.8455	1.7292
SSSF	A_{11}	0.0159	0.1105	0.2136	0.2743	0.3065	0.3247
	A_{12}	0.0116	0.0926	0.2107	0.3089	0.3768	0.4218
	A_{13}	15.2944	5.3778	3.3720	2.5797	2.1637	1.9087
	A_{22}	0.0181	0.1085	0.1924	0.2425	0.2713	0.2887

Table 2. The numerical values of nonlinear natural frequencies of isotropic levy plates at various aspect ratios

Boundary condition	Mode	Nonlinear natural frequency, (rad/sec)					
		Aspect ratio, β					
		$\beta = 0.25$	$\beta = 0.50$	$\beta = 0.75$	$\beta = 1.0$	$\beta = 1.25$	$\beta = 1.50$
SCSC	ω_{11}	363.497	95.2656	45.9213	29.0854	23.2362	22.5013
	ω_{12}	473.004	139.173	80.9642	62.6300	50.0136	42.5369
	ω_{13}	1013.780	258.712	118.987	70.3803	56.2073	60.022
	ω_{22}	2049.33	517.789	234.221	135.156	90.7409	71.5289
SCSS	ω_{11}	254.198	69.3414	35.3535	23.7843	19.653	19.2091
	ω_{12}	366.991	118.946	75.0828	60.2854	42.7581	33.7225
	ω_{13}	838.929	215.700	100.381	60.6915	54.8403	53.8604
	ω_{22}	2058.11	520.456	235.745	136.163	90.1307	66.3526
SSSS	ω_{11}	167.815	49.357	27.454	20.5135	19.3654	19.5027
	ω_{12}	315.027	106.266	69.7313	54.209	40.8385	33.850
	ω_{13}	713.114	185.143	87.458	57.9053	54.0947	56.3444
	ω_{22}	1862.55	471.88	214.423	124.733	82.6781	59.174
SCSF	ω_{11}	47.7087	19.6398	14.874	14.3069	13.9204	13.471
	ω_{12}	116.163	59.5027	43.8963	35.6418	35.2985	35.6283
	ω_{13}	267.400	78.6609	50.5387	48.853	54.5767 - 0.8225i	53.918 - 5.7295i
	ω_{22}	991.382	257.027	122.055	76.6943	54.5767 + 0.8225i	53.918 + 5.7295i
SSSF	ω_{11}	0 + 10472.4i	0 + 2588.97i	0 + 1531.26i	0 + 1194.23i	0 + 1044.08i	0 + 965.899i
	ω_{12}	42.6735	18.8572	14.1528	12.7282	12.169	11.7224
	ω_{13}	103.498	58.4575	45.4109	33.2682	31.362	33.8029
	ω_{22}	305.628	86.0388	50.3588	47.5902	46.406	45.960

From Table 2, it was observed that the nonlinear natural frequencies for the levy plates under investigation exhibit soft-spring type with aspect ratios, except at the third modes of SCSC and SSSS rectangular plates, which deviated from the norm at aspect ratios 1.25 and 1.5 respectively. Furthermore, it was observed that the SCSF rectangular plate boundary conditions yielded positive real-valued complex conjugate nonlinear natural frequencies at the third and fourth modes in the vicinity of aspect ratios of 1.25 and 1.5 respectively. This exposition significantly doubts the stability of the SCSF rectangular plate boundary conditions at third and fourth modes at aspect ratio greater or equal to 1.25 based on the existence of high positive value of the real part of the complex conjugates. However, at aspect ratios less than or equal to 1.0, the SCSF rectangular plate boundary conditions yielded distinct positive nonlinear natural frequencies at the same third and fourth modes respectively. This exposition signifies that within the vicinity of the identified aspect ratios, the SCSF rectangular plate is stable at the third and fourth modes respectively. Also, from Table 2, it was observed that the SSSF rectangular plate boundary conditions yielded positive pure imaginary nonlinear natural frequencies at the first mode for all the selected aspect ratios, which means that the SSSF plate boundary conditions exhibits critical behaviour at the first mode, otherwise it is said to be merely stable. This exposition signifies that in the geometrically nonlinear dynamic analysis of SSSF rectangular plate boundary conditions using modal combination approach, the fundamental nonlinear frequency is sought at second mode instead of the first mode vis-à-vis the linear fundamental frequency of the plate boundary conditions.

Table 3 presents the comparison of the nonlinear fundamental frequencies of the present study's results with the results found in literature. However, the restriction to the use of only the nonlinear fundamental frequencies for the comparison was due to the non-availability of the results for higher mode nonlinear natural frequencies in almost all the literature reviewed.

Table 3. Comparison of Nonlinear fundamental frequencies with results found in literature ($\beta = 1$)

Plate properties	Boundary condition	Investigator	Fundamental frequency (rad/s)	Percentage error
$E = 10.92 \text{ MPa};$ $\rho = 100 \text{ kg/m}^3$ $\mu = 0.3$ $a = 1.0 \text{ m}$ $b = 1.0 \text{ m}$ $t = 0.01 \text{ m}$	SCSC	Present study	29.0854	
		[20]	29.1873	0.349
		[3]	29.2315	0.500
		[11]	29.2627	0.606
	SSSS	Present study	20.5135	
		[20]	20.0990	-2.062
		[3]	20.0916	-2.100
		[11]	20.0818	-2.150
	SCSS	Present study	23.7843	
		[20]	24.0308	1.026
		[3]	24.0113	0.943
		[11]	23.9717	0.782
$E = 30 \text{ GPa}$ $\rho = 25 \text{ kg/m}^3$ $\mu = 0.3$ $a = 10 \text{ m}$ $b = 10 \text{ m}$ $t = 0.2 \text{ m}$	SSSS	Present study	422.785	
		[23]	421.3570	-0.339
	SCSF	Present study	293.666	
		[23]	295.1590	0.506
	SSSF	Present study	261.982	
		[23]	261.2294	-0.288

From Table 3, it was observed from columns 4 and 5 that the present study's results exhibit lower-bound with respect to the results found in literature for the SCSC, SCSS and SCSF rectangular plates respectively. But in the case of SSSS and SSSF rectangular plates, the present study's results exhibit upper-bound with respect to the results found in literature. Also, from column 5 of Table 3, the calculated absolute mean percentage error is 0.186, which shows that the present study's results are in good agreement with the results found in literature.

4. Conclusions

This paper evaluates the nonlinear natural frequencies for Levy plates using weak-form variational principle in algebraic orthogonal polynomial displacement functions. The nonlinear fundamental frequencies evaluated in this work are in good agreement with the results found in literature that were evaluated using other approximation techniques such as the finite element methods. Therefore, it is here-upon concluded that the efficacy of weak-form variational principle technique in dynamic analysis of Levy plates is very satisfactory vis-à-vis other approximation techniques. Also, it is here-unto concluded that the use of algebraic polynomial shape functions in vibration analysis of Levy plates yields satisfactory approximations to all the boundary conditions of the plates. Furthermore, it was observed in this work that the nonlinear natural frequencies for SCSF rectangular plate boundary conditions in the third and fourth modes exhibited positive real valued complex conjugates at aspect ratios 1.25 and 1.50 respectively. Thus, it is here concluded that at aspect ratio greater than or equal to 1.25 in the third and fourth modes respectively, the motions of vibration of SCSF rectangular plate are not stable under geometrically nonlinear conditions. Also, it was observed that the nonlinear natural frequencies in the first mode at various aspect ratios for SSSF rectangular plate exhibited positive pure imaginary values. The exposition here shows that at first mode, the nonlinear motion of vibration for SSSF rectangular plate under modal combination analysis is pole centred, otherwise it is said to exhibit critical behaviour or it is merely stable at the first mode. Therefore, it is here concluded that using modal combination analysis, the nonlinear fundamental frequency for SSSF rectangular plate is ought to be sought in the second mode instead of the usual first mode.

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